



## RESEARCH ARTICLE

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## Aerosol Dynamics in the Near Field of the SCoPEX Stratospheric Balloon Experiment

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**Key Points:**

- Improved models of stratospheric solar geoengineering require accurate predictions of aerosol formation in aircraft wakes
- A new 3D aerosol microphysics model predicts distribution and optical density of sulfate and calcite aerosols in the SCoPEX propeller wake
- This model is a step toward quantitative comparison between models and measurements of aerosol formation in stratospheric wakes

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**Abstract** Stratospheric aerosol injection (SAI) might alleviate some climate risks associated with accumulating greenhouse gases. Reduction of specific process uncertainties relevant to the distribution of aerosol in a turbulent stratospheric wake is necessary to support informed decisions about aircraft deployment of this technology. To predict aerosol size distributions, we apply microphysical parameterizations of nucleation, condensation, and coagulation to simulate an aerosol plume generated via injection of calcite powder or sulfate into a stratospheric wake with velocity and turbulence simulated by a three-dimensional (3D) fluid dynamic calculation. We apply the model to predict the aerosol distribution that would be generated by a propeller wake in the Stratospheric Controlled Perturbation Experiment (SCoPEX). We find that injecting  $0.1 \text{ g s}^{-1}$  calcite aerosol produces a nearly monodisperse plume and that at the same injection rate, condensable sulfate aerosol forms particles with average radii of  $0.1 \text{ }\mu\text{m}$  at 3 km downstream. We test the sensitivity of plume aerosol composition, size, and optical depth to the mass injection rate and injection location. Aerosol size distribution depends more strongly on injection rate than injection configuration. Comparing plume properties with specifications of a typical photometer, we find that plumes could be detected optically as the payload flies under the plume. These findings test the relevance of in situ sampling of aerosol properties by the SCoPEX outdoor experiment to enable quantitative tests of microphysics in a stratospheric plume. Our findings provide a basis for developing predictive models of SAI using aerosols formed in stratospheric aircraft wakes.

### 1. Introduction

Solar geoengineering, or solar radiation modification (SRM), is the largescale intentional manipulation of Earth's radiative budget for the purpose of offsetting radiative forcing introduced by accumulating greenhouse gases. The introduction of aerosol or aerosol precursors to create an artificial reflective stratospheric aerosol layer, known as stratospheric aerosol injection (SAI), was identified as a promising form of SRM in the National Research Council (2015) report. Current literature suggests that, coupled with emissions cuts, SAI could reduce increases in global mean temperature and extreme climate events (droughts and extreme precipitation) associated with increased greenhouse gas burdens. Recent works, including P. J. Irvine et al. (2016a, 2017, 2019) and Kravitz et al. (2013) detail these climatic effects and their global distributions.

SAI is contingent on the ability to artificially generate stratospheric aerosol. Several methods of transporting aerosol into the stratosphere have been considered in the literature, including artillery shells, hoses, balloons and aircraft (Battisti et al., 2009; Davidson et al., 2012). Of these, the release of aerosol or aerosol precursor from aircraft is considered both feasible and relatively inexpensive (P. J. Irvine et al., 2016b), therefore the remainder of this work will focus on understanding SAI in the context of aircraft injection. While many studies have examined stratospheric SRM in general circulation models, only two studies (Benduhn et al., 2016; Pierce et al., 2010) have simulated the creation of an aerosol plume from an aircraft. Both studies used zero-dimensional aerosol microphysics driven by the Turco and Yu (1997) plume dilution study. Observational data on aerosol behavior in stratospheric aircraft plumes is similarly limited (Fahey et al., 1995, 1993; Turco & Yu, 1997). More accurate predictions of the SAI aerosol distribution produced by aircraft will require better models that are tested and improved by more SAI-specific observations.

To effectively model the climate impacts of SAI, plume development from realistic injection scenarios must first be modeled, validated, and applied as initial conditions for global circulation models (GCMs). While the majority of current GCM studies represent SAI as a uniform solar reduction (Kravitz et al., 2011, 2013),

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a subset of models explicitly simulate aerosol microphysics, allowing size distributions to freely evolve (Dai et al., 2018; Simpson et al., 2019; Tilmes et al., 2009). However, even models with explicit aerosol microphysics implicitly assume that injected aerosol or aerosol precursors are instantaneously mixed into model grid boxes with horizontal scales of 100 km or greater. This approach neglects the near-field response, meaning the small-scale chemical and physical processes which would occur in aerosol plumes generated by aircraft. In the near field, the high-concentration gradients which define a plume will drive nonlinear microphysical processes and chemistry (Cameron et al., 2013). In some cases, near-field microphysics and chemistry will play an important role in determining the aerosol size distribution that mixes into the background stratosphere.

Three options for the generation of stratospheric aerosols have been proposed: (a) the release of gas phase  $\text{SO}_2$  (b) the injection of  $\text{H}_2\text{SO}_4$  gas to form accumulation mode (AM)  $\text{H}_2\text{SO}_4$  aerosol (Pierce et al., 2010) or, (c) the injection of solid aerosols (e.g., calcite, diamond, alumina, and titania) (Ferraro et al., 2011; Pope et al., 2012; Teller et al., 1996; D. K. Weisenstein et al., 2015).

The radiative efficiency, or the radiative forcing per injected aerosol mass, of sulfate aerosol is strongly size dependent. Backwards scattering efficiency decreases beyond an optimal radius of 150 nm, while gravitational particle settling increases with particle size. Heckendorn et al. (2009) found that injection of gas phase  $\text{SO}_2$ , which slowly oxidizes to form  $\text{SO}_3(\text{g})$  before producing sulfate aerosol, tends to form particles with radii substantially larger than the optimal size ( $r > 1 \mu\text{m}$ ). This tendency serves to decrease the effective radiative cooling per unit mass of injected sulfur. Direct injection of AM- $\text{H}_2\text{SO}_4$ , which immediately forms sulfate aerosol, to better control, or “steer,” aerosol size distribution was proposed by Pierce et al. (2010). Only two studies, Pierce et al. (2010) and Benduhn et al. (2016), have looked at the near-field dynamics of the injection of AM- $\text{H}_2\text{SO}_4$  into an aircraft wake. Each of these used expanding plume models with nucleation and coagulation parameterizations for the  $\text{H}_2\text{O}$ - $\text{H}_2\text{SO}_4$  system which remain largely unvalidated for the high density, high temperature conditions of an aircraft wake. In both cases, a simulated high-density injection of AM- $\text{H}_2\text{SO}_4$  was allowed to develop until adequately diluted into pre-existing background particles, at which point the final sulfate aerosol size distribution was extracted.

Solid aerosols for SAI were first proposed by Teller et al. (1996), and a growing literature has looked at the use of solid engineered particles. Several studies have found that for the same radiative forcing, solid aerosols can produce less ozone loss, stratospheric heating and forward scattering as compared to sulfate aerosols (Ferraro et al., 2015; Keith et al., 2016; Pope et al., 2012; D. K. Weisenstein et al., 2015). However, no studies have investigated behavior of a high-density injection of solid particles on spatial scales relevant to aircraft injection.

GCM and two-dimensional (2D) global scale studies have investigated the impacts of introducing AM- $\text{H}_2\text{SO}_4$  as compared to  $\text{SO}_2(\text{g})$ . Findings of these works suggest that SAI using AM- $\text{H}_2\text{SO}_4$  may give better control of particle size and produce higher radiative efficiency relative to an  $\text{SO}_2(\text{g})$  injection (Dai et al., 2018; Vattioni et al., 2019; D. Weisenstein et al., 2019). In the case of solid engineered particles, a global 2D chemical transport model has been used to predict climactic outcomes of SAI using injection of monodisperse solid aerosols (Keith et al., 2016; D. K. Weisenstein et al., 2015). All of these results, however, depend on the size distribution of the aerosols injected into the grid-boxes of a global model, and in the case of solid aerosol or AM- $\text{H}_2\text{SO}_4$  that size distribution will be strongly dependent on plume-scale behavior.

Observations of stratospheric aerosol formation via gaseous  $\text{SO}_2$  injection from volcanic eruptions and aircraft emissions provide information about the behavior of aerosols in the stratosphere and the largescale climatic impacts of point source injections (Friedl, 1997; Haywood et al., 2013; Proctor et al., 2018; Rahmes et al., 1998; Robock, 2000; Trenberth & Dai, 2007). However, data from volcanic eruptions do not strongly constrain estimates of response to continuous or nonpoint source injection and aircraft emissions exhaust sulfate in low concentrations compared to what would be expected if implementing SAI. In both cases, high concentrations of coemitted constituents detract from the usefulness of the data, and neither observation relates to solid engineered particles. Therefore, these observations alone cannot effectively validate an SAI-like injection. Measurements in the exhaust plume of the Concorde aircraft highlighted uncertainties in the microphysics of the  $\text{H}_2\text{O}$ - $\text{H}_2\text{SO}_4$  aerosol system in an aircraft wake. Classical condensation and coagulation parameterizations were unable to replicate particle growth in the exhaust plume to achieve observed size

distributions, and particle characteristics were highly sensitive to the homogeneous nucleation rate (Brown et al., 1996; Fahey et al., 1995; Kärcher et al., 1995). This prompted the theory that accounting for ion-mediated nucleation was necessary to accurately predict particle growth under high temperature conditions, although this has not been further validated (Yu & Turco, 1997). Data sets from convective events in the lower stratosphere and volcanic eruptions give insight into mid-field (10–1,000 km) to global transport of aerosols, but do not provide information about near-field plume and particle dynamics. No direct observations of stratospheric AM–H<sub>2</sub>SO<sub>4</sub> or solid aerosol exist to provide avenues for model validation.

Robust understanding of the microphysical processes of these materials will require the intentional codevelopment of model and observational experiments focused on constraining uncertainties of aerosol nucleation, condensation, evaporation and coagulation under SAI-like conditions in the stratosphere.

In this study, we investigate aerosol growth dynamics using a 3D microphysical model driven by the velocity and turbulence fields from high-resolution computational fluid dynamics (CFD) simulations. We base our simulation on the design of the Stratospheric Controlled Perturbation Experiment (SCoPEX) which was proposed by Dykema et al. (2014) to investigate aerosol microphysics in a turbulent wake and to collect information about catalytic chemistry and dynamics of the lower stratosphere (~20 km altitude). SCoPEX will use propellers to steer the payload and generate a region of turbulence into which aerosol or aerosol precursor will be injected. The payload will then navigate downstream to sample the plume at multiple locations.

We incorporate these design features in our model to study the injection of a solid aerosol and vapor-phase aerosol precursor from a balloon payload. We vary initial aerosol mass fluxes and examine the size distributions and optical detectability of produced plumes. We are particularly interested in the ability to produce a plume where the size distribution is representative of a radiatively efficient SAI deployment from an aircraft. Model inputs may be updated by the SCoPEX team with measured plume dynamical parameters to compare model and observations. In the concluding discussion, we speculate about how this model may be applied to study SAI in aircraft plumes.

## 2. Model

Our modeling system is made up of three sub-models: (a) a commercial CFD package to compute a 3D time invariant distribution of velocity and turbulent diffusion; (b) a MATLAB-based advection-diffusion scheme to carry aerosol through the CFD-derived velocity-diffusion fields; and (c) a sectional aerosol microphysical model that computes the evolution of the aerosol size distribution in each grid cell via nucleation, condensation and coagulation. All modeling apart from the CFD analysis used MATLAB 2017a, and all computations were run on the FASRC Cannon cluster.

### 2.1. CFD With ANSYS Fluent

The CFD domain is constructed with ANSYS DesignModeler and meshed with ANSYS Meshing, both in the ANSYS Workbench package 17.2. Following the best practice guideline for CFD (Franke et al., 2004), the inlet and outlet of the domain should be at least  $5d$  upstream and  $15d$  downstream of the propellers, where  $d$  is the propeller diameter (1.9 m). The inlet is 20 m upstream and the outlet is 3,000 m downstream, satisfying both requirements. The entire CFD domain is a cylinder of diameter 200 m and length 3 km. The domain consists of 12 million cells in an adaptive mesh, with the finest cell (0.025 m) in the vicinity and downstream of the propellers. The maximum cell expansion ratio is 1.2 to ensure no sudden change of the mesh resolution.

The simulation is conducted with the finite-volume solver ANSYS Fluent. We use a steady Reynolds-averaged Navier-Stokes (RANS) solver with the standard  $k$ - $\epsilon$  turbulence closure scheme. The Semi-Implicit Method for Pressure-Linked Equations (SIMPLE) algorithm is used for pressure-velocity coupling. The Least Squares Cell Based method is used for discretization of gradients, the Second Order scheme is used for discretization of pressure, and the Second Order Upwind scheme is used for discretization of momentum, turbulent kinetic energy and turbulent dissipation rate. The boundary conditions are: constant velocity of  $1 \text{ ms}^{-1}$  at the inlet, uniform reference (gauge) pressure of 0 Pa at the outlet, and free-slip at the outer surface. The propellers are modeled as 3D Fan Zones with 1.8 Pa pressure jump. The far-field ambient pressure is set

at 5529 Pa and the fluid density is  $0.0889 \text{ kg m}^{-3}$ . These conditions simulate two propellers moving forward steadily at  $1 \text{ m s}^{-1}$  at an altitude of  $\sim 20 \text{ km}$ .

## 2.2. Advection and Diffusion Model

Advection and diffusion are driven by the steady state velocity fields produced by Fluent. Velocities and diffusion coefficients are interpolated from the unstructured Fluent grid to the model grid using a Delaunay triangulation of the scattered sample points via MATLAB's scattered interpolant function, using a linear method. The algorithm is detailed by Amidror (2002).

The coefficients for turbulent diffusion are computed from turbulent viscosity using the standard ANSYS Schmidt number of 0.7. As described in the section above, we assume a stratosphere with laminar background flow of  $1 \text{ ms}^{-1}$ , and do not attempt to represent high intensity turbulence. All turbulence accounted for in this model arises from motion of the propellers. This omission seems reasonable given the estimated waiting time of  $\sim 24 \text{ h}$  between turbulent encounters for a typical stratospheric air parcel (Vanneste, 2004). Balloon experiments, like SCoPEX, can likely avoid significant turbulence if isolated from lee waves, jet stream adjustments, and vigorous atmospheric convection.

To balance both the constraints required to achieve numerical stability and the need to resolve a small enough spatial scale to mimic injection from a nozzle, the desired domain of 3 km was broken into three nested grids. This allowed for high resolution at the nozzle "inlets" while keeping the simulation computationally tractable. Measured from the propeller, these regions are defined as: Box 1, from  $-3$  to 6 m, with 0.1 m resolution, Box 2 from 6 to 100 m at 1 m resolution, and Box 3 from 100 to 3,000 m with 3 m resolution. These boxes encompass the entire x dimension of the Fluent domain. Box boundaries are based on the average magnitude of the X-direction velocity, leveraging lower velocity values to increase the numerically stable spatial step, as detailed below.

We apply a second order accurate Lax-Wendroff advection scheme with a Superbee flux limiter to prevent spurious oscillations in locations with high velocity gradients while also preventing excessive numerical diffusion. 3D versions of this method were adapted from LeVeque (2002), and employed using logic outlined by Smaoui and Radi (2001). Operator splitting enabled the separate computation of diffusion across each time step. This was done using a simple 3D finite difference scheme as described by Hyman et al. (1997). The time steps for each domain were chosen to ensure that the Courant-Friedrichs-Lewy (CFL) value in each dimension did not exceed 0.8, and was, on average, substantially smaller. This value was chosen based on testing in Box 1, where we found a total domain mass error of 0.39% as compared to simulation with a maximum CFL equal to 0.065. This corresponded to timesteps of 0.008, 0.08, and 0.9 for Box 1, 2, and 3, respectively. The maximum diffusion coefficients of a given box were at most two orders of magnitude smaller than maximum velocity values in the same domain, such that achieving advection stability requirements ensured numerical stability of the diffusion scheme.

## 2.3. Sectional Aerosol Model

The nucleation, condensation and coagulation subroutines of the Atmospheric and Environmental Research (AER) 2-D chemical transport model (Weisenstein et al., 1997, 2006) were translated into MATLAB and coupled to the advection and diffusion code described above. They were applied, as discussed below, to separately analyze the continuous release of calcite ( $\text{CaCO}_3$  (s)), or sulfuric acid ( $\text{H}_2\text{SO}_4$  (g)), into the modeled propeller wake.

This model uses 40 logarithmically spaced sulfate aerosol bins, ranging in size from  $3.9 \times 10^{-4}$  to  $3.2 \mu\text{m}$ , representing a doubling of particle volume from bin to bin. Solid calcite aerosol is represented by 8 sectional bins ranging from 0.28 to  $2.8 \mu\text{m}$ , in which the number of monomers per particle is doubled between adjacent bins. Simulated sulfate aerosols are formed through homogeneous nucleation of  $\text{H}_2\text{SO}_4$  and  $\text{H}_2\text{O}$  vapor, with explicit simulation of condensation, evaporation and coagulation (Weisenstein et al., 1997, 2006). Calcite aerosols are injected as monomers, and size distributions are modified only by coagulation among particles. This study did not include simultaneous injection of gaseous and solid species, and therefore only homogeneous coagulation is considered.

Because temperature, pressure, and relative humidity for this experiment are held constant, sulfate aerosol weight fraction is based on the total gas phase concentration of sulfuric acid at a given location. Aerosol composition is used to calculate the nucleation rate and critical cluster radius as presented in Vehkamäki (2002). High-density injection of  $\text{H}_2\text{SO}_4(\text{g})$  was found to exceed the validity range of the Vehkamäki nucleation rate parameterization of  $10^{10} \text{ cm}^{-3} \text{ s}^{-1}$  in some spatial locations. To account for this, the AER nucleation module was modified to identify locations where the concentration of sulfuric acid exceeds the threshold of barrierless kinetic particle formation, and to calculate the nucleation rate via the parameterizations of Määttänen et al. (2018). Following the precedent of Pierce et al. (2010), all sulfuric acid vapor is forced to condense if the nucleation rate exceeds the total particle density in a given spatial location.

The condensation and evaporation scheme, as given by Hamill and Yue (1979), and applied in the AER 2D model, are applied without modification. These processes are grouped into a subroutine which calculates the growth or loss of mass from a sulfate aerosol. The change in particle volume per second per surface area is determined from the sulfate aerosol weight fraction and surface vapor pressure, where a curvature correction for the Kelvin effect is included in the calculation (D. K. Weisenstein et al., 1997). The coagulation schemes derived and applied for sulfate and solid aerosol by D. K. Weisenstein et al. (1997, 2006, 2015) have been applied without modification.

### 3. Results

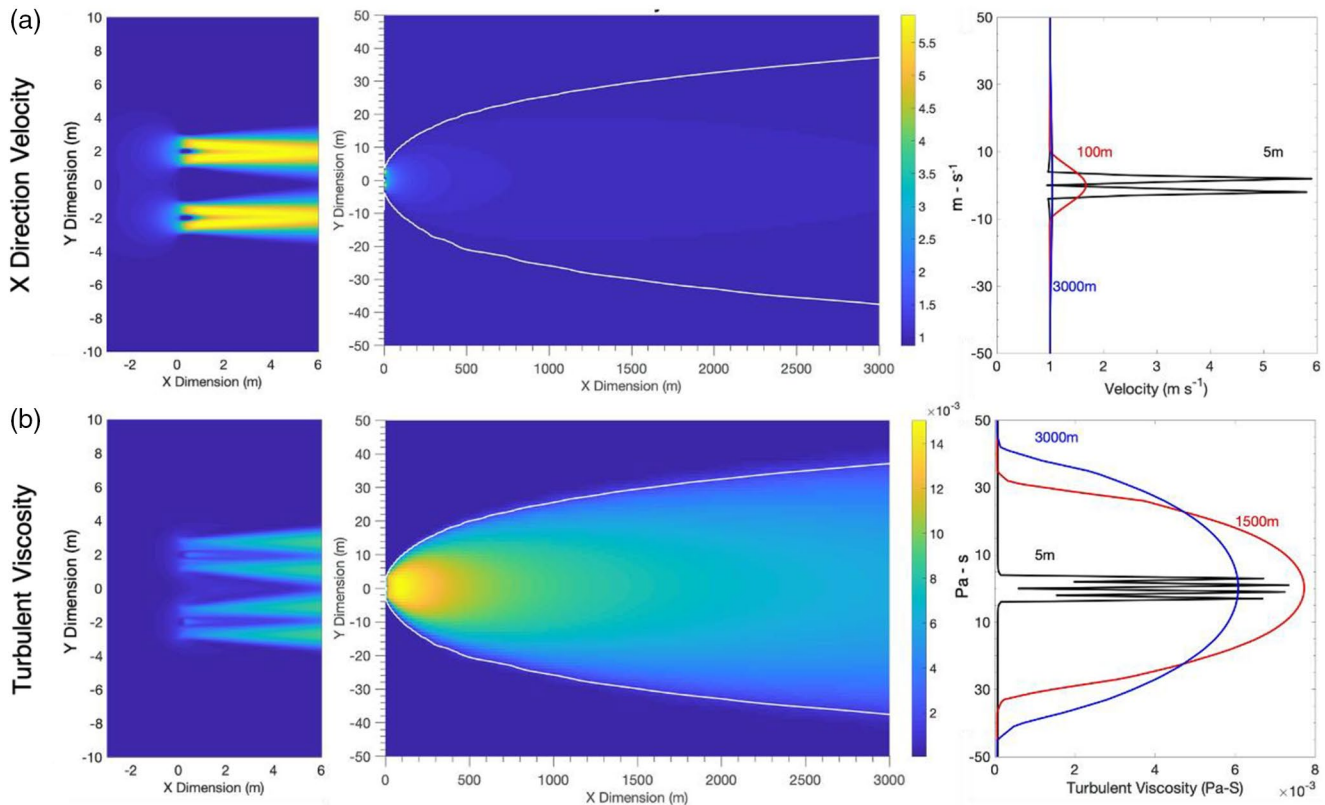
The advection-microphysics model configuration described above is applied to understand the near-field dynamics of calcite ( $\text{CaCO}_3$ ) and sulfuric acid aerosol proposed for use in SCoPEX. As we explore in the concluding discussion, this problem is similar to the aerosol release into a stratospheric jet wake.

For the SCoPEX mission to provide observations relevant to SRM, it needs to produce downstream aerosols with radii within the range of roughly 0.2–1.0  $\mu\text{m}$ . For calcite, the objective is to maintain a high fraction of the aerosol in monomer form, while for sulfate an ideal distribution would have a peak volume mean diameter (VMD) of 0.6  $\mu\text{m}$  (Dykema et al., 2016). The least desirable outcome of the SCoPEX mission is to create particles substantially larger than these ideal sizes, as both particle lifetime and radiative efficiency decrease with increasing particle size. Generation of largely smaller than ideal particles, while imperfect for assessing radiative efficiency relevant to SAI, does not serve to increase particle sedimentation rates within the plume. Distributions centered on small particle sizes may continue to evolve beyond the domain of the study.

The velocity and turbulent viscosity fields from the CFD simulation are shown in Figure 1. These fields form the basis of the simulation environment. We find peaks in x-direction velocity,  $u$ , directly downstream from the modeled propeller centers, with an absolute maximum value of  $6.33 \text{ m s}^{-1}$ . By 1,500 m downstream from the inlet locations, the velocity is reduced to the imposed background flow of  $1 \text{ m s}^{-1}$ . Turbulent viscosity, used as a measure of particle mixing with background air, exhibits a narrow distribution of maximum values  $\sim 10 \text{ m}$  downstream from simulated propellers. With increasing distance downstream, the spatial distribution of turbulent viscosity widens, attaining a full width at half maximum (FWHM) of 60 m by 1,500 m downstream. Notably missing from these fields is the wake of the balloon itself, which is assumed to be sufficiently far from the payload to avoid wake crossing. In additions, this simulation assumes a laminar stratospheric background flow, neglecting the potential impacts of breaking gravity waves. Results may be altered by initializing the CFD simulation with a higher turbulent ratio.

The first challenge to achieving an accurate near-field delineation is the numerical representation of the physical injection nozzles. Currently, no literature provides technical specifications relevant to aerosolizing material for SAI. Leveraging ongoing laboratory studies, we assume that for SCoPEX, precipitated  $\text{CaCO}_3$  powder with roughly monodisperse size distribution centered at  $\sim 0.5 \mu\text{m}$  VMD will be aerosolized using the expansion of powder suspended in high pressure  $\text{CO}_2$  through a 1–2 mm nozzle. Sulfuric acid vapor will be emitted from a heated reservoir through a nozzle of a few millimeter diameter. The millimeter-scale injector orifices are too small to directly resolve, instead we model injection as a 3D gaussian distribution of mass flux into the model grid, where the size of that distribution represents the scale of the nozzle spray pattern  $\sim 10 \text{ cm}$  from the orifice. We use gaussian distributions with a FWHM of 0.3 and 1.0 m for the calcite and  $\text{H}_2\text{SO}_4$ , respectively.



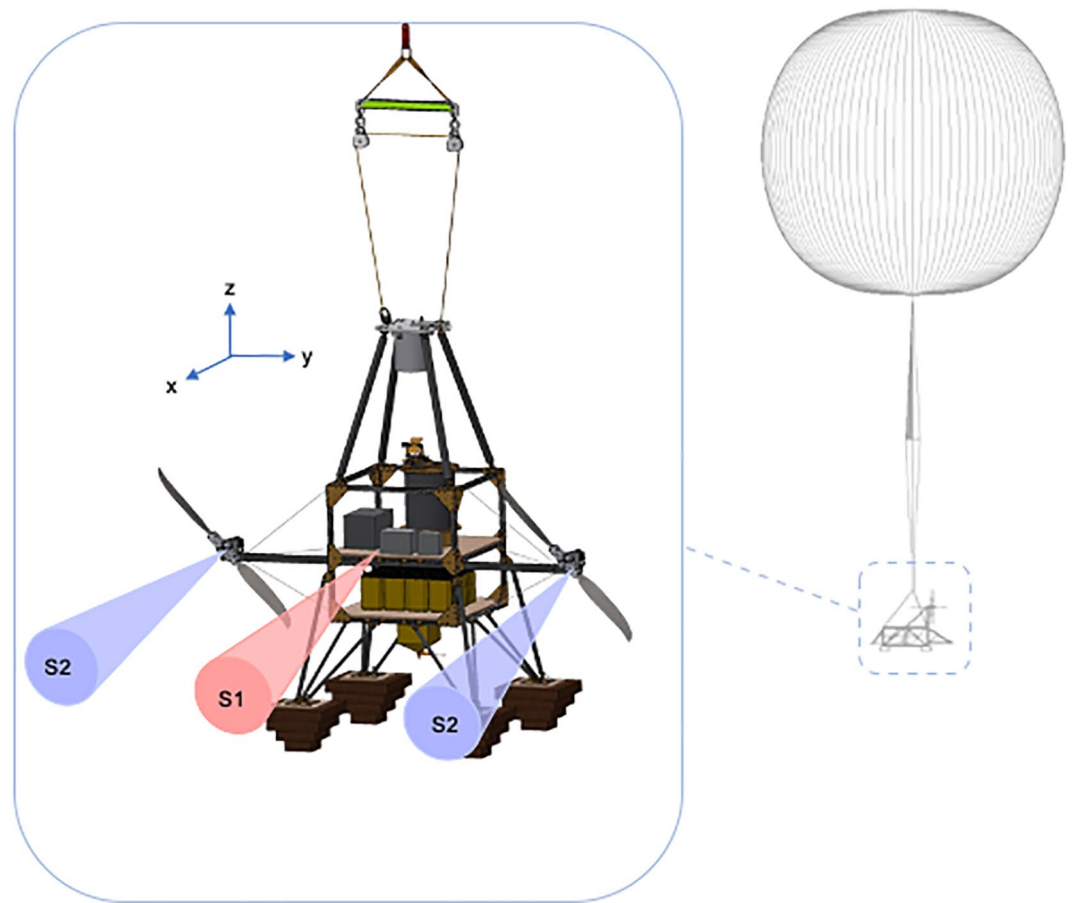


**Figure 1.** Shown above are the steady state x-direction velocity,  $\mu$ , and turbulent viscosity fields generated by ANSYS Fluent. Left panels show the genesis of disruptions to background x direction flow of  $1 \text{ m s}^{-1}$  at the inlets to two propellers, respectively located, at  $(0,2)$  and  $(0,-2)$  meters. The center panel shows the entire domain, from 0 to 3 km, where the imposed white line contours  $1 \text{ m s}^{-1}$  in plot (a), and contours 10% of the absolute maximum turbulent viscosity in plot (b). Note Y direction scaling differs between the center and left panels. The right panel shows cross sections of velocity (a) and turbulent viscosity (b) through the Y plane at varying X locations.

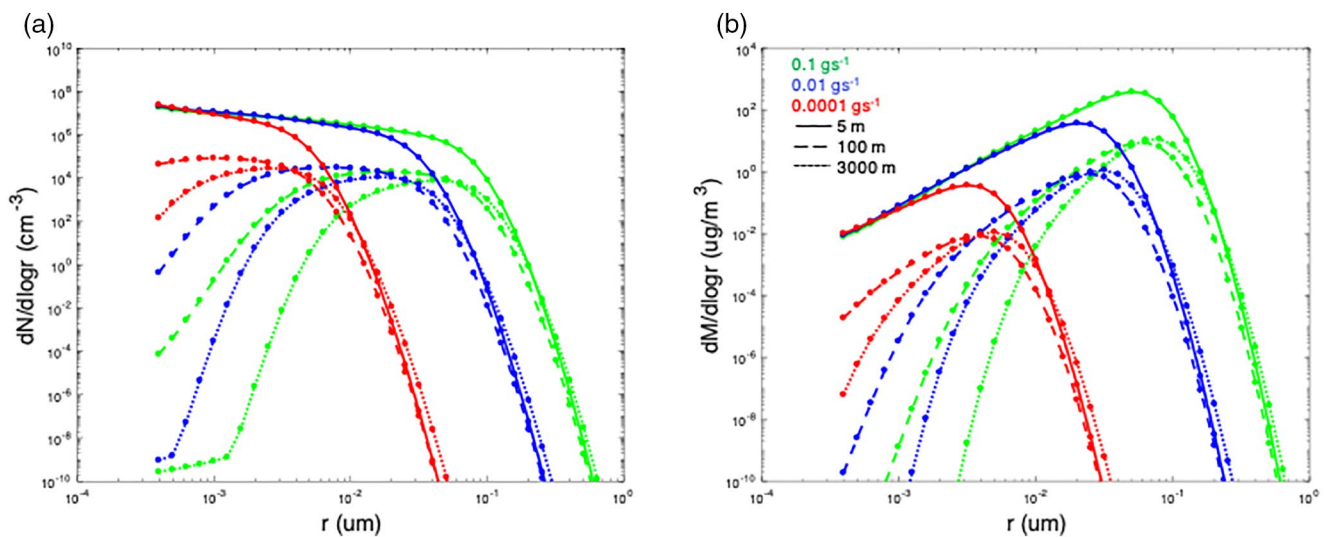
The SCoPEX experiment team considered two configurations for the continuous injection of material from the payload (Figure 2): scenario 1 (S1), a single point injection between the propellers; and scenario 2 (S2), injection from the center of each propeller. While scenario 1 is the simplest to implement, we were not sure if it would adequately mix the material into a sufficiently large volume. Defining the plume “edge” as locations where the total particle number density drops below  $1 \text{ cm}^{-3}$ , we find plume diameter at 3 km to be insensitive to injection scenario for injection of both  $\text{AM-H}_2\text{SO}_4$  and calcite. For sulfate, the FWHM (y-dimension diameter) of the plume is consistent across all scenarios and mass injection rates, with a value of 93 m. This suggests that injection at or between the propellers does not serve to greatly alter the characteristics of the particles’ experienced velocity field. Calcite injection yields similar results, with an average plume FWHM of 70 m and a standard deviation of 3 m, or one spatial grid cell.

These findings support that scenario 1 is effective at dispersing material into the propeller wake. We also find that the final size distribution of particles within the plume is impacted by injection location. Comparison between S1 and S2 across the same mass injection rate shows an average percent reduction in VMD at 3,000 m of 11% and 16% for sulfate and calcite injection, respectively. The final size distribution of particles within the plume, while sensitive to location of injection, is more dependent on mass injection rate, as discussed in the next section. For this reason, the remainder of our results will focus on outcomes of an S1 injection.

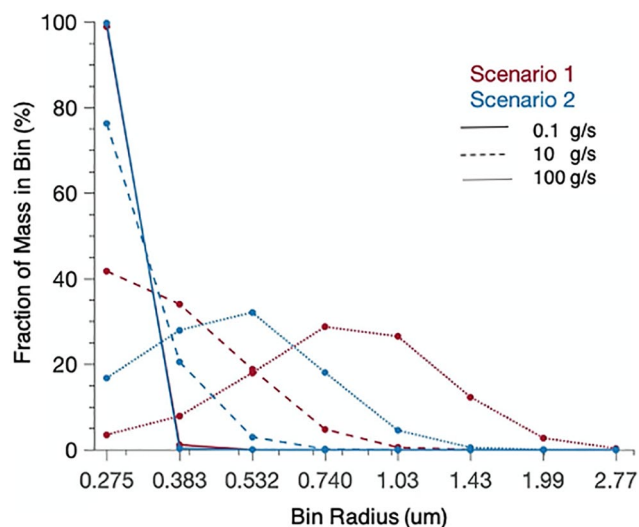
Mass injection rates were varied by factors of 10 and 1,000 to test the influence of initial particle number density on the final plume aerosol size distribution. These results (Figures 3–5) show that particle size increases with injection rate and downstream location. Increasing the mass flux of  $\text{H}_2\text{SO}_4$  from  $0.0001$  to  $0.1 \text{ g s}^{-1}$ , reduces the negative skew of the downstream (3 km) size distribution by decreasing the share of smaller



**Figure 2.** A simplified schematic of the proposed SCoPEX balloon and payload system. Potential injections scenarios 1 and 2 are shown in red and blue, respectively, as injections from the center of the system (S1), or the center of each propeller hub (S2).

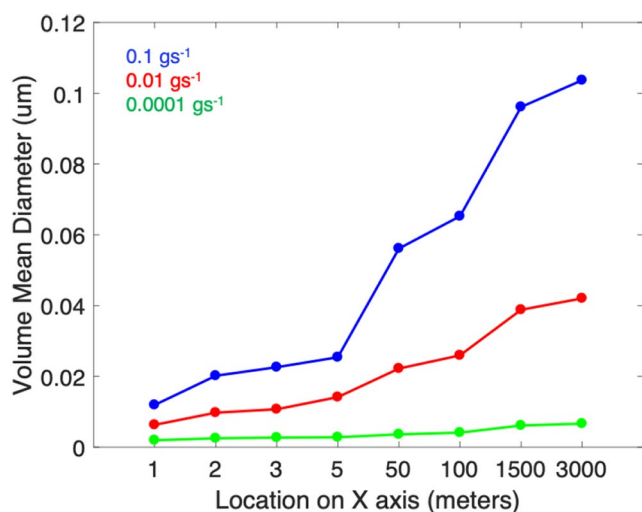


**Figure 3.** Steady state sulfate aerosol number (a) and mass (b) size distributions predicted by the AER model. Red, blue and green lines represent injection rates of 0.1, 0.01, and 0.0001  $\text{g s}^{-1}$ , under scenario 1, respectively. Solid, dashed and dotted lines represent average integrated aerosol number density and mass through the plume at sampling locations downstream of 5, 100, and 3000 m, respectively. AER, Atmospheric and Environmental Research.



**Figure 4.** Calcite aerosol size distributions, shown as the fraction of total mass of in each sectional bin where the  $x$ -axis markers represent the central radius of each sectional size bin. These distributions represent the percent of total aerosol mass in the final 100 m of the plume across the full domain. Results are shown for three injection rates, 0.1, 10, and 100  $\text{g s}^{-1}$ , for injection scenarios 1 (red) and 2 (blue).

rate are not driving changes in particle growth. After 6 m, virtually no gas phase  $\text{H}_2\text{SO}_4$  remains in the system, such that further changes to particle size can only be due to particle coagulation. In Figure 5, the rate of change of VMD of the plume has significantly slowed by 1,500 m, implying that the plume is disperse enough that coagulation between particles is no longer relevant. This finding tells us that to generate a sulfate plume with particle VMD of 0.6  $\mu\text{m}$ , we must increase the initial mass injection rate; simply allowing the plume more time develop is unlikely to produce the desired result.



**Figure 5.** Steady state volume-mean particle diameter integrated over the  $y$ - $z$  plane at increasing distances downstream from injection location. Results are shown for condensable  $\text{H}_2\text{SO}_4$  injected via scenario 1 at rates 0.1, 0.01, and 0.0001  $\text{g s}^{-1}$  as blue, red and green solid lines, respectively. In all cases a monotonic increase in average particle size is observed with distance downstream.

particles and shifting the distribution peak toward radii of  $\sim 0.1 \mu\text{m}$ . Shifting calcite injection from 0.1 to 100  $\text{g s}^{-1}$  similarly reduces the share of monomer particles in the final 100 m of the plume to favor multi-monomer ( $r > 0.275 \mu\text{m}$ ) aggregates (Figure 4). The differences in the resulting VMD of low versus high initial mass fluxes, increase with distance downstream, or longer time in the system (Figure 5).

While findings for both materials follow the same trends, we rely on coagulation to promote particle growth when injecting gas phase  $\text{H}_2\text{SO}_4$ , whereas the objective of injecting 0.55  $\mu\text{m}$  monomer calcite is to minimize coagulation between particles; to conserve monomer. Increasing initial mass flux of  $\text{H}_2\text{SO}_4$  by a factor of 1,000, from 0.0001 to 0.1  $\text{g s}^{-1}$ , increases the number average particle radius at 3 km by 1,500%, from 0.0025, to 0.04  $\mu\text{m}$ , a desirable shift. Comparatively, the lowest calcite flux of 0.1  $\text{g s}^{-1}$ , is the most desirable, maintaining 99% of the total mass in the final 100 m of the plume in monomer form. Increasing mass injection rate to 10 and 100  $\text{g s}^{-1}$ , with an S1 injection, shifts peak mass loading to favor particles of radii 0.5 and 0.75  $\mu\text{m}$ , respectively.

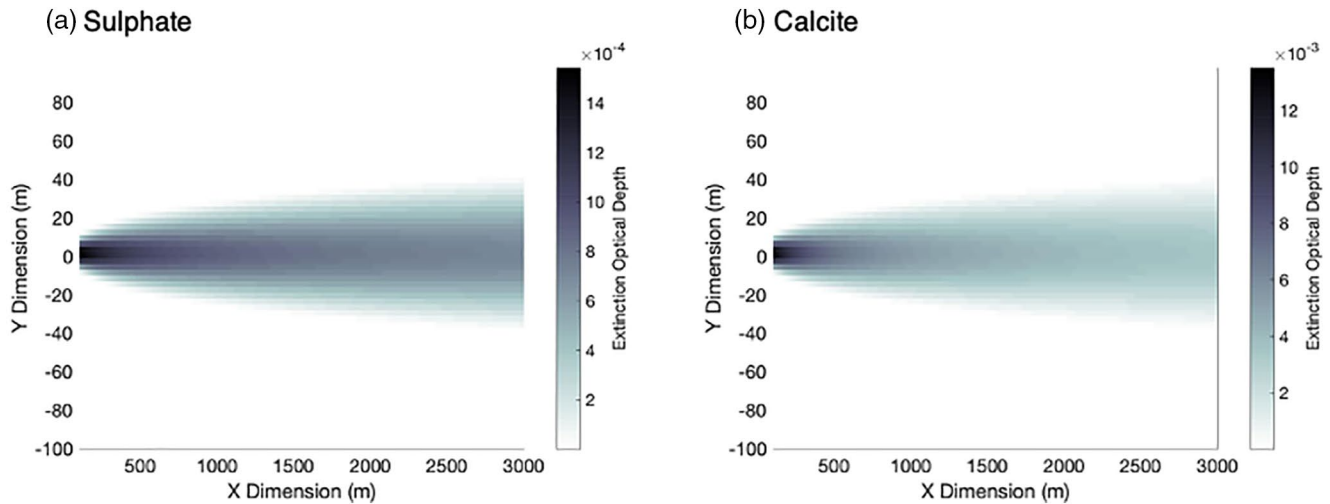
All tested mass injection rates display an increase in particle VMD with distance downstream. For sulfate, this results in particles with VMD of 0.006, 0.04, and 0.1  $\mu\text{m}$  for injection rates of 0.0001, 0.01, and 0.1  $\text{g s}^{-1}$ , respectively. This trend corroborates the notion that coagulation processes drive the final particle size distribution (Pierce et al., 2010). The relative differences in VMD between injection rates are the lowest closest to the nozzle, where nucleation dominates, suggesting changes to nucleation

### 3.1. Optical Detection

Steady state number density distributions across the final domain were leveraged to investigate detectability of the plume. Using Mie scattering theory, the extinction optical depth was calculated by vertically integrating down columns in the  $y$ - $z$  plane. Figure 6 shows the relative optical thickness of the sulfate and calcite aerosol plumes formed via scenario 1 with an injection rate of 0.1  $\text{g s}^{-1}$ . Calcite exhibits greater optical thickness by an order of magnitude at 550 nm, with an average value of  $8.6 \times 10^{-4}$  and maximum of 0.014 across the domain, as compared to sulfate, with an average of  $9.4 \times 10^{-5}$  and maximum of 0.001.

To better understand detectability of a theoretical plume against background stratospheric aerosol we extract a scattering measurement at 550 nm as could apply to observations from a narrowband, narrow field view photometer with azimuthal/zenith pointing capability. We assume an altitude of 21 km and solar elevation angle of  $60^\circ$  for a SCoPEX-like flight occurring at solar noon at  $40^\circ\text{N}$  during stratospheric turnaround. The observing instrument is assumed to be situated on the payload gondola (Figure 1), 200 m from the edge of the plume and positioned to observe the plume 1 km downstream of the termination of a scenario 1 type injection of calcite aerosol.





**Figure 6.** Extinction optical depth integrated vertically through all columns in the plume from 100 to 3,000 m. Plots a and b show results for  $0.1 \text{ g s}^{-1}$  injections of condensable  $\text{H}_2\text{SO}_4$  and calcite, respectively. The resulting number density of calcite aerosol is  $490 \text{ cm}^{-3}$  on the centerline at a downstream distance of 1,000 m, predominantly as monomers. At the same location, sulfate aerosol is present at a density of  $2 \times 10^5 \text{ cm}^{-3}$ . Aerosol optical depths were derived from Mie scattering theory at 550 nm, using refractive indices for sulfate and calcite stated in Dykema et al. (2016).

Under these conditions, a notional photometer can employ established viewing geometries for measuring diffuse solar radiation (Torres et al., 2014). We assume that the gondola is physically located so it can scan the principal plane by altering its zenith pointing angle. By locating itself such that the plume is located at a zenith pointing angle of  $20^\circ$ , the ratio of the aerosol plume scattering to the background is maximized due to angular dependence of the corresponding scattering phase functions. The scattered radiation received by the photometer is proportional to product of the angular scattering coefficient (units of  $\text{m}^{-1} \text{ sr}^{-1}$ ), integrated along the line of sight, and the photometer field of view (in sr).

For an injection rate of  $0.1 \text{ g s}^{-1}$  and a vertical displacement of 200 m between the plume and the gondola, the ratio of the angular calcite scattering to the combined Rayleigh (molecular) and aerosol background angular scattering is about 10, when integrated over the line of sight. Assuming a photometer with a field of view comparable to that of Murphy et al. (2016) spanning a half-angle of  $2.5^\circ$  or  $0.0065 \text{ sr}$ , the expected scattered optical power incident on the photometer is about 15 nW. This level of optical power is about 5 orders of magnitude above the noise floor for a 1 ms integration time for compact, high-reliability optical detectors working at 550 nm. However, for this optical power the maximum achievable SNR is about 6,000, due to photon shot noise. While a real instrument will achieve a reduced SNR due to optical and electrical inefficiencies, we expect adequate SNR to confidently detect the plume with a fast-scanning radiometer via the solar radiation it scatters.

#### 4. Discussion and Conclusions

We found injection scenarios that allow the SCoPEX balloon experiment to generate sulfate or calcite aerosol plumes with size distributions that are relevant to understanding aerosol deployment in stratospheric aircraft wakes. Injection of gaseous  $\text{H}_2\text{SO}_4$  with a mass injection rate of  $0.1 \text{ g s}^{-1}$  using a single injection location forms particles with a VMD of  $0.1 \mu\text{m}$ , which are smaller than the ideal but large enough to be experimentally relevant. Injection of solid calcite at  $0.1 \text{ g s}^{-1}$  produces an almost perfectly monodisperse distribution.

Injection rate and injected material largely dictate the aerosol size distribution and optical thickness of the resulting plume. Plume diameter is largely determined by the fluid dynamics of the propeller wake and is insensitive to aerosol injection rate and location. In all cases, increasing mass injection rate increased the average optical thickness and the peak of particle size distribution of the plume. If injecting condensable  $\text{H}_2\text{SO}_4$ , larger mass injection rates favor the formation of accumulation mode particles, while low-injection

rates decrease total number density, limiting condensation and coagulation as mechanisms to grow particles. Alternatively, increasing mass injection rate of calcite tends to shift particle size distributions to favor aerosols consisting of two or more monomer units. Under the same injection strategy, a calcite plume exhibits significantly larger per mass scattering efficiency as compared to a plume generated via AM-H<sub>2</sub>SO<sub>4</sub>.

We find that the aerosol plumes in the SCoPEX wake could be detected by broad-band sun-angle-resolving photometer located just below the wake. The primary measurement of aerosol properties in SCoPEX will be by in situ aerosol instruments. Our scattering calculations suggest that these observations could be complimented by a small photometer mounted on the SCoPEX gondola as it is maneuvered under the wake. This detection method may also be relevant to studies of aerosol formation in stratospheric aircraft plumes.

There are two important limitations to our study. First, our model incorporates Brownian coagulation, but omits gravitational and turbulent coagulation. In the latter case, particle inertia, impacting relative particle velocities, and wind field shear, would serve to increase particle collisions. These transport effects would likely dominate when the Stokes number is close to 1, and may only be a small concern given our background flow of 1 m s<sup>-1</sup> (Benduhn, 2008; Riemer et al., 2007). For particle distributions less than 1 μm, Brownian coagulation largely dominates regardless of atmospheric turbulence, suggesting that these omissions are unlikely to have influenced outcomes of sulfate injection in which particle sizes largely remain below 1 μm. As particles become larger, the inertia of an individual particle within a flow field becomes significant and turbulent coagulation becomes relevant. Therefore, calcite distributions with particles larger than micrometer size are likely lacking coagulation kernel contributions from turbulent processes. Second, the calcite coagulation sticking coefficient was taken at a value of 1 but could have a real value anywhere between 0 and 1. These omissions and assumptions need to be tested and verified through iterations with observations and laboratory studies.

This work provides a foundation to use observations from SCoPEX or similar experiments to quantitatively evaluate models of aerosol microphysics using observations. This modeling framework can be used in combination with (a) aerodynamic measurements, (b) lidar observations of plume extent, and, (c) in situ particle size observations to constrain microphysical quantities such as the coagulation kernel. The overall goal would be to test models of aerosol microphysics by identifying weaknesses in the parameterizations of high-density aerosol nucleation, condensation and coagulation in a stratospheric wake. With adequate validation, this model can then be extended to model solar geoengineering deployment technologies from an aircraft. Such a study would need to extend nucleation parameterizations to include ion-induced nucleation and interaction of aerosol particles and precursors with other species from the effluent streams in the aircraft wake. This would naturally build on the work of Pierce et al. (2010) and Benduhn et al. (2016), in the context of condensable H<sub>2</sub>SO<sub>4</sub>, and extend findings to solid aerosols, like calcite.

Iteration between model development and observations of stratospheric aerosol formation will be needed to provide more accurate near-field predictions of aerosol properties used to initialize global models of stratospheric response to aerosol solar geoengineering.

## Data Availability Statement

Data generated and investigated in this work are openly available on Harvard Dataverse (<https://doi.org/10.7910/DVN/JMHKFT>).

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